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Charged dilatonic black holes Phase transitions and holographic duals

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A 15' presentation on:

- Dilatonic Einstein-Maxwell theory in AdS
- Instability of RN-AdS and phase transitions
- Charged dilatonic black holes
- AdS/CFT: The dual field theory
- Electrical conductivity in a strongly coupled field theory

Our model: Dilatonic Einstein-Maxwell in AdS

$$\mathcal{L} = R - \frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi - V(\phi) - \frac{f(\phi)}{4}F^2$$

Nonminimal coupling between the dilaton and the Maxwell field

$$V(\phi) = -\frac{6}{L^2} + \frac{m^2}{2}\phi^2 + \mathcal{O}(\phi^3), \qquad f(\phi) = 1 + \frac{\alpha}{2}\phi^2 + \mathcal{O}(\phi^3)$$

- **RN-AdS** is solution
- Quite well motivated model

Is RN-AdS the only static BH solution?

Instability of RN-AdS.

• Scalar linear perturbations:

$$\phi = \sum_{l,m} R(r) Y_{lm}(\theta, \varphi) e^{-i\omega t}$$

• Can be recast into a Schroedinger-like equation:

$$\frac{d^2\Psi(z)}{dz^2} + \left[\omega^2 - V_l(r) - m_{\text{eff}}^2(r)\right]\Psi(z) = 0$$



Instability of RN-AdS. Naïve argument

• In AdS a scalar perturbation produces an instability if

$$m_{
m eff}^2 \ll -rac{9}{4L^2}$$
 — Breitenlohner-Freedman bound



Instability below a critical temperature

$$T \ll T_c(\alpha,m^2)$$

also confirmed by

analytical argumentsnumerical integration

What is the endpoint of the instability??

Numerical resolution. Planar symmetry

$$ds^{2} = -g(r)e^{-\chi(r)}dt^{2} + \frac{dr^{2}}{g(r)} + r^{2}(dx^{2} + dy^{2})$$

- 4 fields (2 metric functions, 1 Maxwell field, 1 dilaton)
- 4 independent eqs. (2 Einstein's, 1 Maxwell's, 1 scalar eq.)



New solutions: Charged dilatonic BHs in AdS

Numerical resolution imposing suitable boundary conditions



RN-AdS will decay into this when T<T

Fine, that was nice but...

Universe is ~ flat and at most Λ>0, so...

<u>Why are you studying</u> <u> $\Lambda < 0 \ gravity??</u></u>$

Because Maldacena's conjecture!

J:hams

AdS/Condense Matter dictionary:

Metric tensor



Black hole temperature

Conserved energy-stress tensor

Temperature of the theory

Scalar field:

$$\phi \rightarrow \frac{\mathcal{O}_-}{r^{\lambda_-}} + \frac{\mathcal{O}_+}{r^{\lambda_+}}$$



Expectation values of scalar operators

Gauge field:

 $A_0 \to \mu - \frac{\rho}{r}$

Chemical potential, charge density



An example: High-Tc Superconductors

Electrical Conductivity



It is just gravity!!! (No electrons, No Cooper's pairs...)

/ - Frequency gap
 - Infinite DC conductivity
 - Charged condensate

[Hartnoll, Horowitz, Hertog, Gubser, Roberts, Polchinski, Kachru...] (2008-now)

Results: second order phase transition



How can the condensation of a neutral operator change the transport proprieties of a strongly coupled theory?

Electrical conductivity

Transport phenomena in the dual theory can be studied by **linear perturbations in the bulk**

$$A_{\mu} = [A_0(r), 0, 0, 0] + [0, 0, A_x(r), 0]e^{-i\omega t}$$

$$A_x'' + \left(\frac{g'}{g} - \frac{\chi'}{2} + \frac{f'(\phi)}{f(\phi)}\phi'\right)A_x' + \left[\left(\frac{\omega^2}{g^2} - \frac{A_0'^2 f(\phi)}{g}\right)e^{\chi}\right]A_x = 0$$

In the AdS/CFT dictionary the electrical conductivity can be computed by the asymptotical behavior of the perturbations

$$A_x \sim A_x^{(0)} + \frac{A_x^{(1)}}{r}$$

$$\sigma = -irac{A_x^{(1)}}{\omega A_x^{(0)}} = rac{1-\mathcal{R}}{1+\mathcal{R}}$$
AdS/CFT result

Numerical results Conductivity:



- Negative minimum of $V(z) \rightarrow$ slope of the conductivity at low frequency

- No bound states
- Reminiscent of "Drude peak" in real materials!

Schroedinger potential:

Electrical conductivity

Numerical results:







- Universal behavior at high frequency
- Minimum at low frequency
- Slope in the DC limit



Conclusions

- We study a model in which RN-AdS is unstable
- Second order phase transition toward new charged dilatonic
 BHs below a critical temperature

- Phase transition: neutral scalar condensation
- Unexpected effects on transport proprieties of a strongly coupled field theory
- Charged dilatonic black holes share some proprieties with real materials

Something to think about:

We still have to learn a lot from gravity/field theories dualities

Thanks!



Backup slides

$$T_c = \gamma^{\frac{1}{4}} \frac{1}{8\pi L} \sqrt{\frac{Q}{L}} \left[12 - \frac{1}{\gamma} + k \frac{L}{Q} \sqrt{\frac{2}{\gamma}} \right]$$

$$\left[\frac{2}{\gamma}\right] \qquad \gamma = \frac{\alpha}{\frac{9}{4} + \beta}$$

Num. Sol. of the Field Eqs. Planar symmetry

$$ds^{2} = -g(r)e^{-\chi(r)}dt^{2} + \frac{dr^{2}}{g(r)} + r^{2}(dx^{2} + dy^{2})$$

• Two nonlinear, coupled differential equations:

Scalar equation:

$$\phi''(r) + \left(\frac{g'}{g} - \frac{\chi'}{2} + \frac{2}{r}\right)\phi'(r) - \frac{V'(\phi)}{g} + \frac{A_0'^2 e^{\chi} f'(\phi)}{2g} = 0$$

Einstein equation:

$$\frac{{\phi'}^2}{4} + \frac{{A'_0}^2 e^{\chi} f(\phi)}{4g} + \frac{g'}{rg} + \frac{1}{r^2} + \frac{V(\phi)}{2g} = 0$$

The gauge field and one metric function can be directly integrated

Preliminaries:

- Reisser-Nordstrom BH (in flat spacetime)
- Two horizons:

• The inner horizon is unstable

$$r_{\pm} = M \pm \sqrt{M^2 - Q^2}$$

• Not even an approximate solution of ST

- Here we deal with AdS black holes

- Particles reach infinity in a finite amount of time
- Scalars with negative mass squared (tachyons) can exist if

$$m^2 \ge -\frac{9}{4L^2}$$