



II Workshop on Black Holes

21-22 December 2009 | Instituto Superior Técnico, Lisboa



Instabilities in higher-dimensional spacetimes

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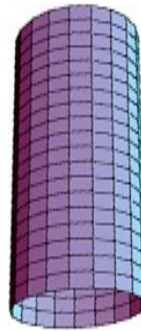
V. C., M. Lemos and M. Marques, Phy. Rev. D, in press (2010).

Instabilities of black objects

The Gregory-Laflamme instability

- **Black string:**

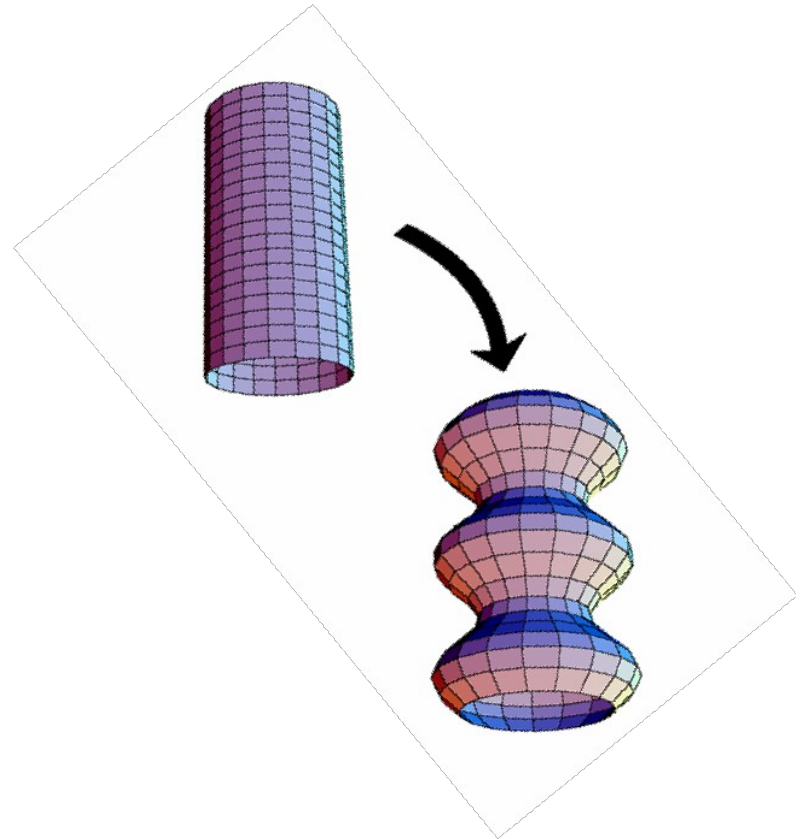
Schwarzschild \times \mathbb{R}



- **Perturb the BS (s-wave)**

String : $S \propto \text{Mass}^2 / L$

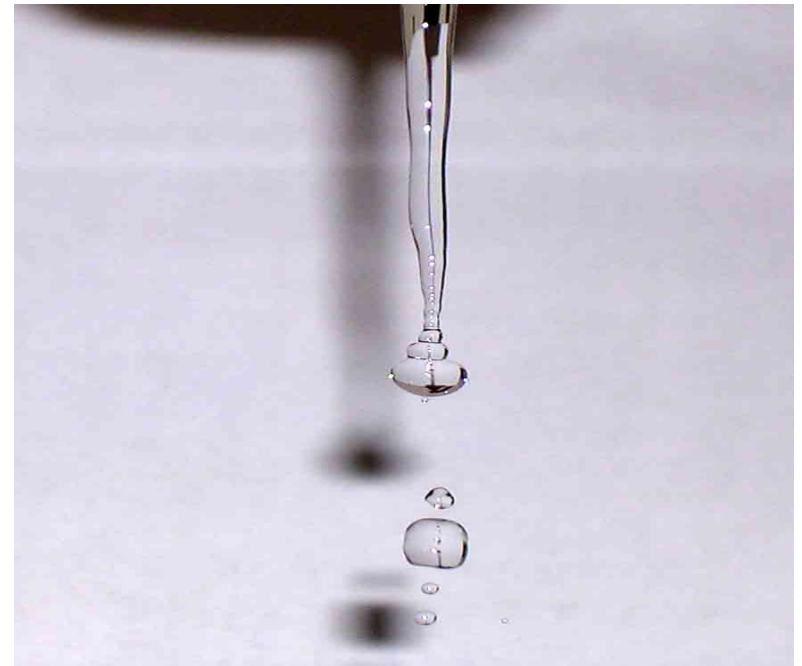
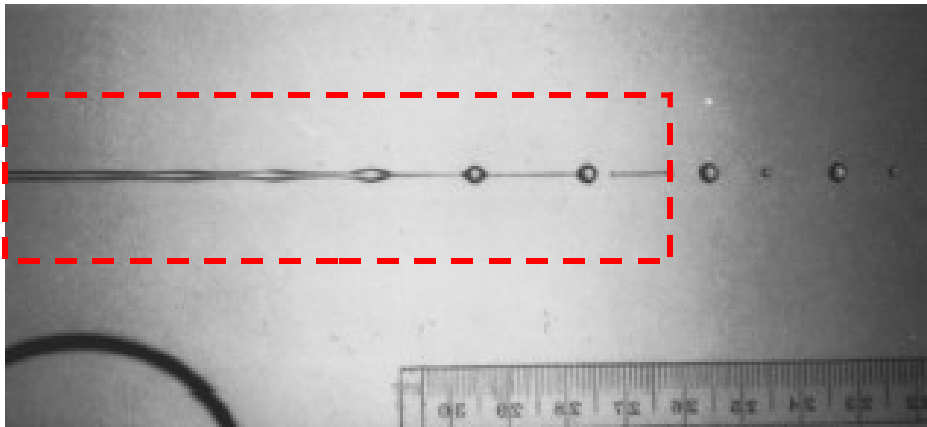
Spherical : $S \propto \text{Mass}^{3/2}$



R. Gregory and R. Laflamme, Phys.Rev.Lett.70:2837,1993

Instabilities of black objects

The Rayleigh-Plateau instability (Plateau 1849, Rayleigh 1878)



$$\text{GL} : Rk_{\text{crit}} \sim \sqrt{D}$$

$$\text{RP} : Rk_{\text{crit}} \sim \sqrt{D}$$

Instabilities of black objects

Ultra-spinning instabilities (R. Monteiro and J. Santos' talks)

Emparan and Myers, JHEP 0309:025,2003; O. J. C. Dias et al, arXiv:0907.2248

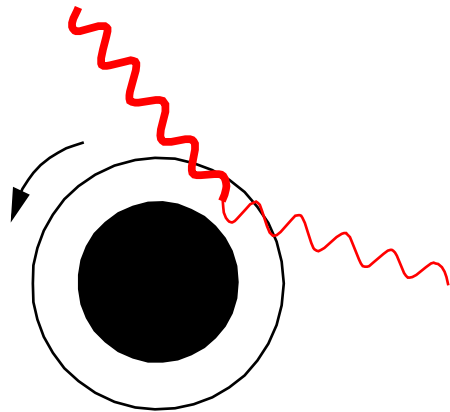
Instabilities of black objects

Super-radiant instabilities (J. Rosa's talk)

Zel'dovich 1971; Damour et al (1976); VC and O.J. C.Dias PRD70:084011,2004.

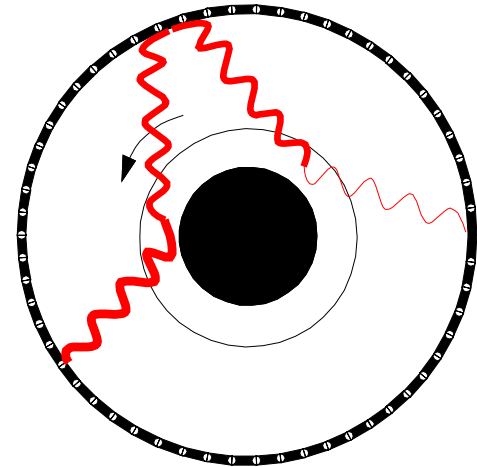
→ **Superradiant scattering
on a rotating BH**

Waves incident upon a Kerr black hole are amplified by superradiant scattering if $\omega < m\Omega$



**Insert a mirror around
the black hole:**

Make a black hole bomb!



(Cardoso, Dias, Lemos, Yoshida, 2004)

AdS acts as box: small Kerr-AdS BHs are unstable (Cardoso and Dias, 2004)

Instabilities of black objects

Instabilities in Einstein-Gauss-Bonnet theories

$$S_{GB} = \int d^D x \sqrt{-g} \frac{\alpha}{(D-3)(D-4)} (R_{abcd}R^{abcd} - 4R_{ab}R^{ab} + R^2)$$

$$ds^2 = -f(r)dt^2 + f(r)^{-1}dr^2 + r^2 d\Omega_{D-2}^2$$

$$f(r) = 1 + \frac{r^2}{2\alpha} - \frac{r^2}{2\alpha} \sqrt{1 + \frac{8\alpha\mu}{r^{D-1}}}$$

Unstable for some range of parameters in D=5,6

Instabilities in Chern-Simons theories (L. Gualtieri's talk)

Instabilities of charged BHs in de Sitter backgrounds

$$ds^2 = -f(r)dt^2 + f(r)^{-1}dr^2 + r^2 d\Omega_n^2$$

$$f(r) = 1 - \lambda r^2 - \frac{2M}{r^{n-1}} + \frac{Q^2}{r^{2n-2}}$$

$$E = q/r^n$$

r_a (Cauchy), $r_b = 1$ (Event), r_c (Cosmological)

$$2M = 1 + Q^2 - \lambda$$

$$\lambda = \frac{r_c^{-4-n} (r_c^{n+2} - r_c^3) (r_c^{n+2} - Q^2 r_c^3)}{r_c^{n+2} - r_c}$$

$$Q_{\text{ext}}^2 = \frac{r_c^n (-2r_c + (n+1)r_c^n - (n-1)r_c^{n+2})}{-r_c (r_c(n+1) - 2nr_c^n + (n-1)r_c^{2n+1})}$$

Instabilities of charged BHs in de Sitter backgrounds

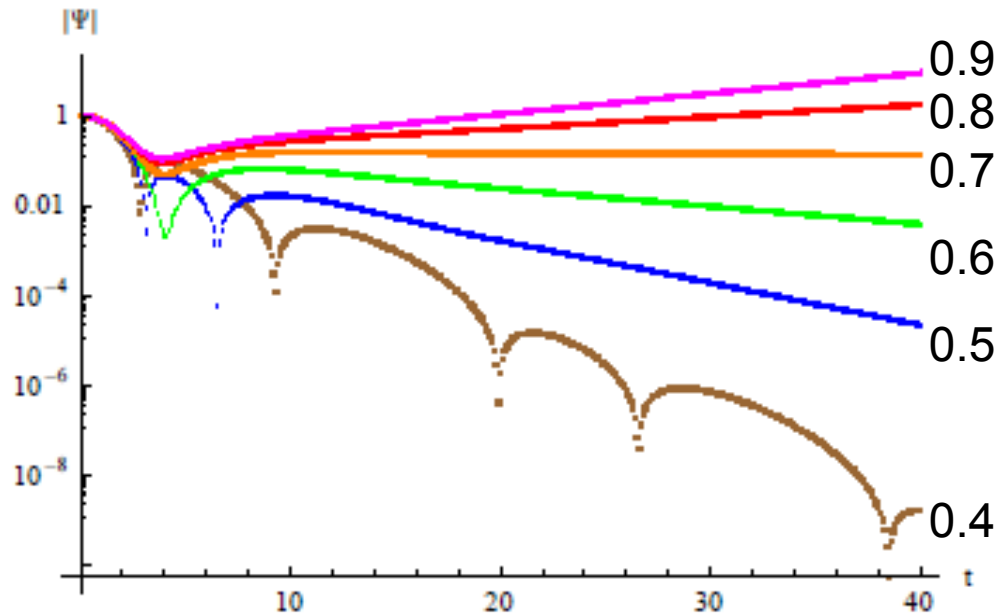
$$\frac{d^2}{dr_*^2} \Psi_{\pm} + (\omega^2 - V_{S\pm}) \Psi_{\pm} = 0, \quad V_{S\pm} = \frac{fU_{\pm}}{64r^2 H_{\pm}^2}, \quad x \equiv \frac{2M}{r^{n-1}}$$

$$\begin{aligned} U_+ = & \left[-4n^3(n+2)(n+1)^2 \delta^2 x^2 - 48n^2(n+1)(n-2) \delta x \right. \\ & \left. - 16(n-2)(n-4) \right] y - \delta^3 n^3 (3n-2)(n+1)^4 (1+m\delta) x^4 \\ & + 4\delta^2 n^2 (n+1)^2 \left\{ (n+1)(3n-2)m\delta + 4n^2 + n - 2 \right\} x^3 \\ & + 4\delta(n+1) \left\{ (n-2)(n-4)(n+1)(m+n^2 K) \delta - 7n^3 + 7n^2 - 14n + 8 \right\} x^2 \\ & + \left\{ 16(n+1) (-4m + 3n^2(n-2)K) \delta - 16(3n-2)(n-2) \right\} x \\ & + 64m + 16n(n+2)K, \end{aligned}$$

$$\begin{aligned} U_- = & \left[-4n^3(n+2)(n+1)^2 (1+m\delta)^2 x^2 + 48n^2(n+1)(n-2)m(1+m\delta)x \right. \\ & \left. - 16(n-2)(n-4)m^2 \right] y - n^3(3n-2)(n+1)^4 \delta (1+m\delta)^3 x^4 \\ & - 4n^2(n+1)^2 (1+m\delta)^2 \left\{ (n+1)(3n-2)m\delta - n^2 \right\} x^3 \\ & + 4(n+1)(1+m\delta) \left\{ m(n-2)(n-4)(n+1)(m+n^2 K) \delta \right. \\ & \quad \left. + 4n(2n^2 - 3n + 4)m + n^2(n-2)(n-4)(n+1)K \right\} x^2 \\ & - 16m \left\{ (n+1)m (-4m + 3n^2(n-2)K) \delta \right. \\ & \quad \left. + 3n(n-4)m + 3n^2(n+1)(n-2)K \right\} x + 64m^3 + 16n(n+2)m^2 K \end{aligned}$$

Instabilities of charged BHs in de Sitter backgrounds

$$D = 11 (n = 9), \quad 1/r_c = 0.8 \quad Q/Q_{\text{ext}} =$$

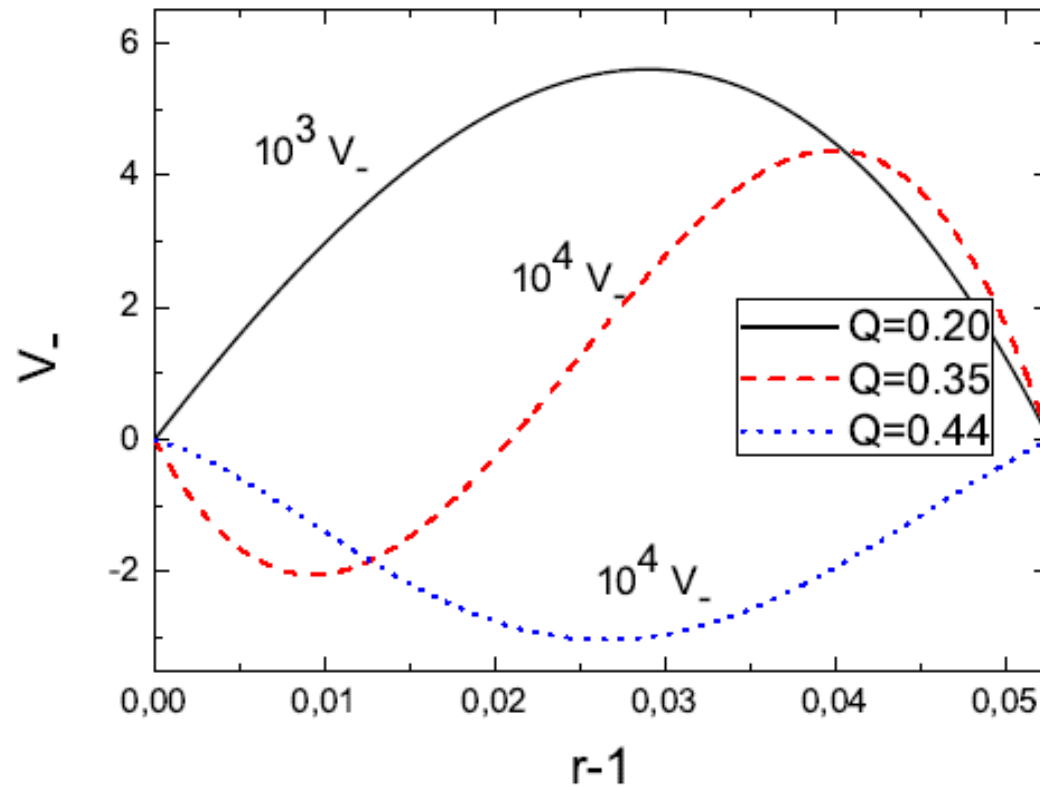


For sufficiently large Q and $D > 6$, BHs are unstable

Konoplya and Zhidenko, Prog.Theor.Phys.111:29-73,2004

Instabilities of charged BHs in de Sitter backgrounds

$$D = 8(n = 6), \quad 1/r_c = 0.95$$



Instabilities of charged BHs in de Sitter backgrounds

$$\frac{d^2\Phi(\omega, r)}{dr_*^2} + \left[\omega^2 - \frac{V_0}{\cosh(\kappa_b r_*)^2} \right] \Phi(\omega, r) = 0$$

$$\omega = \kappa_b \left[- \left(j + \frac{1}{2} \right) i + \sqrt{\frac{V_0}{\kappa_b^2} - \frac{1}{4}} \right], \quad j = 0, 1, \dots$$

	D					
	7	8	9	10	11	$D \rightarrow \infty$
Q/Q_{ext}^N	0.913	0.774	0.683	0.617	0.567	$\sqrt{2/D}$
Q/Q_{ext}^{Num}	0.94	0.78	0.68	0.61	0.55	—

Summary

- Black objects are fascinating from a dynamical viewpoint
- Stability/Instability is an important criterium for relevance
- Stability of charged and/or spinning BHs poorly understood
- Interesting topic for further research

Thank you